lectures on 'hard probes' for quark-gluon plasma diagnosis

- introductory remarks, QGP and hard processes
- lecture 1 heavy quarks and QGP
- lecture 2 jets and jet quenching
- lecture 3 quarkonia

pbm EJC, Villa Clythia, Frejus, Oct. 2013



Quark-gluon plasma –a new state of matterdeconfined quarks and gluons



in relativistic nucleus-nucleus collisions, a new state of matter is produced, in which colored quarks and gluons roam freely

Simulation: UrQMD, Frankfurt

The phase diagram of strongly interacting matter

at low temperature and normal density colored quarks and gluons are bound in colorless hadrons - confinement chiral symmetry is spontaneously broken (generating 99% of proton mass e.g.) 1972 QCD (Gross, Politzer, Wilczek) asymptotic freedom at small distances

<u>at high temperature and/or high density</u> quarks and gluons freed from confinement -> new state of strongly interacting matter 1975 (Collins/Perry and Cabibbo/Parisi)

now called Quark-Gluon Plasma (QGP)



Quark-gluon plasma and the early universe



Evolution of the Early Universe



DARMSTAD

Peter Braun-Munzinger

How to create QGP in the laboratory?



what are 'hard probes'?

Hard probes are observables which involve an energy or mass scale $Q^2 >> (Lambda_QCD)^2$

$$\alpha_{\rm s}(Q^2) = \frac{1}{\beta_0 \ln(Q^2/\Lambda^2)}$$

 $\Lambda \approx 0.1 \text{ GeV for } \alpha_{\rm s}(M_{\rm Z^0} \equiv 91.2 \text{ GeV}) = 0.12$

for such processes, production can be precisely computed in pQCD

in QGP physics, we are interested in studying their propagation in and interaction with the QGP medium



The Large Hadron Collider (LHC)



27 km long, 8 sectors

1232 dipole magnets (15m, 30 tonnes each) to bend the beams Cooled with 120 tonnes of He at 1.9 K

pp: 2808 bunches/ring, each 1.15x10¹¹ protons (8 min filling time) Design luminosity: **10³⁴ cm⁻²s⁻¹**

PbPb: 592 bunches/ring, each 7x10⁷ Pb ions

Design luminosity: 10²⁷ cm⁻²s⁻¹

Transverse r.m.s beam size: 16 µm, r.m.s. bunch length: 7.5 cm

Beam kinetic energy: 362 MJ per beam (1 MJ melts 2 kg copper)

Total stored electromagnetic energy: 8.5 GJ (dipole magnets only)

Overall view of the LHC experiments.





the TPC (Time Projection Chamber) - 3D reconstruction of up to 15 000 tracks of charged particles per event



with 95 m³ the largest TPC ever





560 million read-out pixels! precision better than 500 μm in all 3 dim. 180 space and charge points per track

The interior of the TPC, 2004

first PbPb collisions at LHC at $\sqrt{s} = 2.76$ A TeV

setup for ion collisions: November 4 first collisions with stable beams: November 8 until Dec 6

already in Dec 2010 5 publications in PRL and PLB



Lecture 1

Open Charm and Open Beauty from SPS to LHC Energy

- reminder of basics
- discussion of time scales and open charm conservation equation
- results for RHIC and LHCenergy in pp collisions
- results for RHIC and LHC energy in AA collisions
- an aside on photon production

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FERMIONS matter constituents spin = 1/2, 3/2, 5/2,						
Leptons spin =1/2				Quarks spin =1/2		
Flavor	Mass GeV/c ²	Electric charge		Flavor	Approx. Mass GeV/c ²	Electric charge
VL lightest neutrino*	(0-0.13)×10 ⁻⁹	0		U up	0.002	2/3
e electron	0.000511	-1		d down	0.005	-1/3
𝔑 middle neutrino*	(0.009-0.13)×10 ⁻⁹	0		C charm	1.3	2/3
μ muon	0.106	-1		S strange	0.1	-1/3
𝒫 heaviest neutrino*	(0.04-0.14)×10 ⁻⁹	0		top	173	2/3
τ tau	1.777	-1		b bottom	4.2	-1/3

Remarks on production of open charm and charmonia

- charm quark mass >> Λ_{QCD} production described in QCD perturbation theory
- all calculations employ gluon fusion as starting point
- argument is energy independent until global energy conservation very close to threshold becomes important
- production of charm quark pairs takes place at timescale $1/m_c$ $m_c = 1.5 \text{ GeV} \longrightarrow t_c = 0.13 \text{ fm}$
- to build up wave function of mesons including those with open charm needs about t = 1fm --> charm production and charmed hadron formation are decoupled
- overall cross section is due to production of charm quark pairs
- time scale is much too short to dress the charm quarks essential to take current quarks



Cross section for charm production

based on M. Glueck, J. F. Owens, E. Reya, Phys. Rev. D17 (1978) 2324

in leading order there are 3 important diagrams:





differential cross section

$$\begin{split} \frac{d\sigma^{gg+c\overline{c}}}{dt} &= \frac{\pi\alpha_s^2}{64s^2} \left(12M_{ss} + \frac{16}{3}M_{tt} + \frac{16}{3}M_{uu} \right. \\ &+ 6M_{st} + 6M_{su} - \frac{2}{3}M_{tu} \right) , \quad (A1) \end{split}$$
 with
$$\begin{split} M_{ss} &= \frac{4}{s^2} \left(t - m^2 \right) \left(u - m^2 \right) , \\ M_{tt} &= \frac{-2}{(t - M^2)^2} \left[4m^4 - (t - m^2)(u - m^2) \right. \\ &+ 2m^2(t - m^2) \right] , \end{split}$$
$$\begin{split} M_{uu} &= \frac{-2}{(u - m^2)^2} \left[4m^4 - (u - m^2)(t - m^2) \right. \\ &+ 2m^2(u - m^2) \right] , \quad (A2) \\ &+ 2m^2(u - m^2) \right] , \end{split}$$

$$\begin{split} M_{su} &= \frac{4}{s(u-m^2)} \left[m^4 - u(s+u) \right] \,, \\ M_{tu} &= \frac{-4m^2}{(t-m^2)(u-m^2)} \left[4m^2 + (t-m^2) + (u-m^2) \right] \,, \end{split}$$

total cross section

$$\sigma^{gg + c\overline{c}} = \frac{\pi \alpha_s^2}{64s} \left[12(\frac{2}{3} + \frac{1}{3}\gamma)(1-\gamma)^{1/2} + \frac{16}{3} \left((4+2\gamma) \ln \frac{1+(1-\gamma)^{1/2}}{1-(1-\gamma)^{1/2}} - 4(1+\gamma)(1-\gamma)^{1/2} \right) + 6\left(2\gamma \ln \frac{1+(1-\gamma)^{1/2}}{1-(1-\gamma)^{1/2}} - 4(1+\gamma)(1-\gamma)^{1/2} \right) - \frac{2}{3}2\gamma(1-\gamma) \ln \frac{1+(1-\gamma)^{1/2}}{1-(1-\gamma)^{1/2}} \right]$$

with $\gamma \equiv 4m^2/s \leq 1$.

this result plus NLO/NNLO/FONLL corrections are currently the basis of all open charm calculations (see, e.g., the calculations by Cacciari et al., discussed below.

heavy quark production in pp collisions situation before LHC turned on



Measurement of open charm cross section in pp collisions at the LHC

all LHC experiments contribute

ALICE at pt > 2 GeV/c and 0 < y < 4ATLAS and CMS at pt > 6 GeV/c 0 < y < 2.5LHCb at pt > 2 GeV/c and 2.5 < y < 4

all detectors employ sophisticated Si vertex detectors

D^0 , D^+ and D^{0^*} in 7 TeV pp data



measurements agree well with state of the art pQCD calculations



charm and beauty via semi-leptonic decays

Inclusive electron spectrum from 2 PID methods: TPC-TOF-TRD and TPC-EMCAL



direct γ from pQCD

charm and beauty electrons compared to pQCD



- ALICE data complimentary to ATLAS measurement at higher pt (somewhat larger y-interval)
- good agreement with pQCD
- at upper end of FONLL range for $p_t <$
- 3 GeV/c where charm dominates

arXiv:1205.5423 ATLAS: PLB707 (2012) 438 FONLL: Cacciari et al., arXiv:1205.6344

a first try at the total ccbar cross section in pp collisions



beauty cross section in pp and ppbar collisions



Open heavy flavor production and the QGP

1. m_q >> Lambda_QCD charm quark production is independent of the medium formed in the collision (see above)

2. propagation of heavy quarks in the medium can be used to diagnose it

energy loss – thermalization – hydrodynamic flow

interaction with the hot/dense QCD medium

energy loss

- dependence on medium density and volume
- color charge dependent (Casimir factor) $\rightarrow \Delta E_{gluon} > \Delta E_{gluon}$
- parton mass dependent (dead cone effect: Dokshitzer & Kharzeev, PLB 519(2001)199) $\rightarrow \Delta E_{u,d,s} > \Delta E_{c} > \Delta E_{b}$

thermalization

dependence on transport properties of the medium

Formation time of quarkonia

heavy quark velocity in charmonium rest frame:

v = 0.55 for J/ ψ see, e.g. G.T. Bodwin et al., hep-ph/0611002

minimum formation time: t = radius/v = 0.45 fm

see also: Huefner, Ivanov, Kopeliovich, and Tarasov, Phys. Rev. D62 (2000) 094022; J.P. Blaizot and J.Y. Ollitrault, Phys. Rev. D39 (1989) 232 **formation time of order 1 fm**

formation time is not short compared to plasma formation time especially at high energy

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formation time of open charm hadrons not well understood

presumably similar to charmonia

separation of time scales for initial hard process and late hadronization/hadron formation is called "factorization"

rigorously proven for deep inelastic scattering





no medium effect $\sigma_{c\bar{c}} = 1/2 \left[\sigma_{D^+} + \sigma_{D^-} + \sigma_{D^0} + \sigma_{\bar{D}^0} + \sigma_{\Lambda_c} + \sigma_{\bar{\Lambda}_c} \dots \right]$ medium effects on charmed hadrons affect redistribution

of charm, but not overall cross section

it is not consistent with the charm conservation equation to reduce all charmed hadron masses in the medium for an enhanced cross section



gluon radiation by a quark traversing a medium

from Dokshitzer & Kharzeev, Phys.Lett. B519 (2001) 199-206 we get for the probability of radiation of a gluon with energy ω by a quark with mass M and energy E

$$dP = \frac{\alpha_{\rm s} \ C_F}{\pi} \ \frac{d\omega}{\omega} \ \frac{k_{\perp}^2 \ dk_{\perp}^2}{(k_{\perp}^2 + \omega^2 \theta_0^2)^2}, \qquad \theta_0 \equiv \frac{M}{E}$$
$$k_{\perp}^2 \simeq \sqrt{\hat{q} \ \omega} \qquad \hat{q} \equiv \rho \ \int \frac{d\sigma}{dq^2} \ q^2 \ dq^2 \qquad C_F = \frac{N_c^2 - 1}{2N_c}$$

here the density of scatterers in the medium is encoded in q[^]

'dead cone' effect for charm quarks



Figure 1: Ratio of gluon emission spectra off charm and light quarks for quark momenta $p_\perp=10~{\rm GeV}$ (solid line) and $p_\perp=100~{\rm GeV}$ (dashed); $x=\omega/p_\perp$.

now open charm and open beauty in AA collisions

how to quantify the effect of the medium?

 R_{AA} =yield(AA)/(N_{coll} yield(pp))

 $R_AA = medium/vacuum$

R_AA = 1 if no dense medium is formed or if one looks at electro-weak probes

D meson signals in Pb Pb collisions



suppression of charm at LHC energy

pp reference at 2.76 TeV: measured 7 TeV spectrum scaled with FONLL cross checked with 2.76 TeV measurement (large uncertainty due to limited luminosity)



energy loss for all species of D-mesons within errors equal - not trivial energy loss of central collisions very significant - suppr. factor 5 for 5-15 GeV/c

suppression of charm at LHC energ



Suppression only for Strongly Interacting Hard Probes



photons, Z and W scale with number of binary collisions in PbPb – not affected by medium \rightarrow demonstrates that charged particle suppression is medium effect: energy loss in QGP

charm Quarks also Exhibit Elliptic Flow



model description of energy loss and flow of Dmesons



both are determined by transport properties of the medium (QGP) simultaneous description still a challenge for some models

comparison of suppression for b-quarks and c-quarks



Measurement of the fireball temperature via photon emission

Photons and neutral mesons measured via the conversion method in the ALICE TPC, see, .e.g, M. Wilde (ALICE coll.) QM2012





γ-ray tomography of ALICE



R [cm]

method

- Direct Photon Signal: $\gamma_{direct} = \gamma_{inc} \gamma_{decay} = (1 \frac{\gamma_{decay}}{\gamma_{inc}}) \cdot \gamma_{inc}$
- Double Ratio: $\frac{\gamma_{inc}}{\pi^0} / \frac{\gamma_{decay}}{\pi^0_{param}} \approx \frac{\gamma_{inc}}{\gamma_{decay}}$ if > 1 direct photon signal \rightarrow cancellation of uncertainties
- Numerator: Inclusive γ spectrum per π^0
- Denominator: Sum of all decay photons per π^0 Decay photons are obtained by a cocktail calculation
- Photons and π^0 s are measured via conversion method $\pi^0 \to \gamma + \gamma, \ \gamma \to e^+e^-$

Inclusive photon measurement in Pb-Pb collisions



Final result



average T = 304 +/- 51 MeV

highest ever measured temperature

Interpretation in terms of fireball parameters

Sukanya Mitra, Payal Mohanty, Sabyasachi Ghosh, Sourav Sarkar, Jan-e Alam. Mar 4, 2013.

e-Print: arXiv:1303.0675

TABLE I: The values of various parameters - thermalization time (τ_i) , initial temperature (T_i) and hadronic multiplicity dN/dy - used in the present calculations.

2 76 TeV

VOIVIV	1.10 101
centrality	0-40%
$\frac{dN}{dy}$	1212
$ au_i$	0.1 fm
T_i	$553 { m MeV}$
T_c	$175 { m MeV}$
T_{f}	$100 { m MeV}$
EoS	Lattice QC

SNIN

note: T_i = 3.2 x T_c



Thermal photons also exhibit flow

most hydro-models produce little flow for thermal photons





Figure 4: Inclusive photon $v_2^{\gamma,\text{inc}}$ and decay photon $v_2^{\gamma,\text{bg}}$ in 0–40% Pb-Pb collisions.

Figure 5: Direct-photon $v_2^{\gamma,\text{dir}}$ in 0–40 % Pb-Pb collisions.

O. Linnyk, V.P. Konchakovski, W. Cassing, E.L. Bratkovskaya. Apr 25, 2013.

e-Print: arXiv:1304.7030

Ralf Rapp (Texas A-M, Cyclotron Inst. & Texas A-M). Oct 2011. 30 pp.

Published in Acta Phys.Polon. B42 (2011) 2823-2852

see the 2 papers for a (conflicting) analysis

Phenix results on thermal photons and Rapp's interpretation arXiv:1110.4345



Fig. 14. Transverse-momentum spectra (left panel) and elliptic flow (right panel) of direct photons in 0-20% central Au-Au($\sqrt{s}=200 \text{ AGeV}$) collisions at RHIC, compared to PHENIX data [49, 54]. The curves are calculations [55] with a realistic fireball evolution employing thermal QGP and hadronic rates which are "dual" around T_c , corresponding to Fig. 9.

In this approach, large initial temperatures are not required, and the spectral shapes are due to very strong radial flow. Precision measurements are needed to separate the different approaches.